

Controllability under constraints

Enrique Zuazua

FAU - AvH

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The control problem

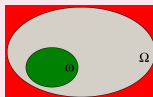
Let $n \geq 1$ and $T > 0$, Ω be a simply connected, bounded domain of \mathbb{R}^n with smooth boundary Γ , $Q = (0, T) \times \Omega$ and $\Sigma = (0, T) \times \Gamma$:

$$\begin{cases} y_t - \Delta y = u1_\omega & \text{in } Q \\ y = 0 & \text{on } \Sigma \\ y(x, 0) = y^0(x) & \text{in } \Omega. \end{cases} \quad (1)$$

1_ω = the characteristic function of ω of Ω where the control is active, $y^0 \in L^2(\Omega)$ and $u \in L^2(Q)$ so that (5) admits an unique solution

$$y \in C([0, T]; L^2(\Omega)) \cap L^2(0, T; H_0^1(\Omega)).$$

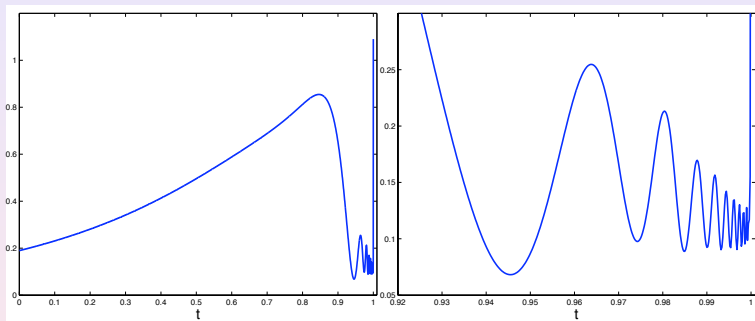
$y = y(x, t) = \text{solution} = \text{state}$, $u = u(x, t) = \text{control}$



Goal: Drive the system to rest:

$$y(x, T) \equiv 0.$$

Numerical simulations, back to Glowinski and Lions in the 80's and 90's



This is so since controls are just restrictions to ω of solutions of the adjoint system

$$-\varphi_t - \Delta\varphi = 0$$

with initial data φ^T at time $t = T$ in very badly conditioned space, the one of the norm one observes.

Bad news

- 1 Controls oscillate dramatically as time approaches the final time.
- 2 This produces oscillations in the state too.
- 3 This effect is further accentuated when the time horizon T is short, as $T \rightarrow 0$.
- 4 This makes the controllability results of little use in many contexts in which the state represents a density (population dynamics, Math Biology).
- 5 This appears systematically in all models of reaction-diffusion type.

Consequences on optimal control

This not only occurs for controllability problems.

Optimal control problems that penalise the final state $\|y(T)\|$ will experience the same behaviour:

$$\min_{u \in L^2(\omega \times (0, T))} \frac{1}{2} \left[\int_0^T \int_{\omega} u^2 dx dt + K \int_{\Omega} |y(x, T)|^2 dx. \right]$$

The reason for this pathology is that, again, the control is, according to the Optimality System (OS), of the form

$$u = -\varphi 1_{\omega}$$

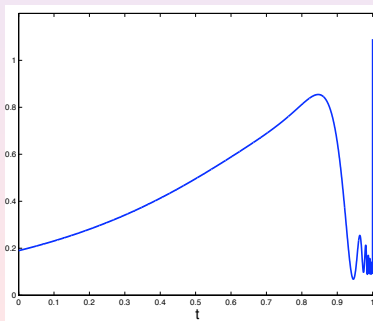
where

$$\begin{cases} \varphi_t + \Delta \varphi = 0 & \text{in } Q \\ \varphi = 0 & \text{on } \Sigma. \end{cases} \quad (2)$$

In the context of null-control, when the goal is to drive the state y to rest ($y(T) \equiv 0$) we know that the final datum φ^T of the adjoint system lies in a very large space that, in terms of the Fourier coefficients on the basis of the eigenfunctions of the Laplacian, can be written as follows:

$$\sum_{j \geq 1} |\hat{\varphi}_j^T|^2 \exp(-c\sqrt{\lambda_j}) < \infty,$$

and this explains the singular behaviour of controls as $t \sim T$:



Can we really expect null-control with nonnegative controls?

The answer is NO! When solving

$$\begin{cases} y_t - \Delta y = u1_\omega & \text{in } Q \\ y = 0 & \text{on } \Sigma \\ y(x, 0) = y^0(x) & \text{in } \Omega. \end{cases} \quad (3)$$

with $y^0 \geq 0$ and $u \geq 0$, then, by the comparison principle of solutions of the heat equation, we have that

$$y \geq z$$

where z is the solution of the heat equation without control. Thus

$$y(x, T) \geq z(x, T) > 0.$$

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And so?

Relax the non-negativity constraint of the control a little bit to

$$u \geq -\delta.$$

And then,

Theorem

For all $\delta > 0$ there exists a positive time-horizon $T(\delta) > 0$ such that the null-control of the equation can be achieved, the control being so that

$$u \geq -\delta.$$

The same result holds for semilinear dissipative heat equations, parabolic equations with variable smooth coefficients, etc.

Sketch of the proof

Goal: Actually prove that

$$\|u\|_{\infty} \leq \delta.$$

Two steps procedure:

- 1 *Phase 1:* Do nothing during time interval $[0, T - 1]$ and let the solution decay.
- 2 *Phase 2:* Once $y(T - 1)$ is small enough, control it to zero in the last time-interval $[T - 1, T]$.

Two well-known key facts

- 1 *Dissipativity.*
- 2 *Controllability.* The controllability of the model, so that in time $[0, 1]$, the system is controllable and with a bound on the cost of control:

$$\|u\|_{L^\infty(\omega \times (0,1))} \leq C^* \|y^0\|_{L^2(\Omega)}.$$

Second proof: Control along a path of steady-states

¹ Allowing to link steady-states, solutions of

$$\begin{cases} -\Delta y = u \mathbf{1}_\omega & \text{in } \Omega \\ y = 0 & \text{on } \partial\Omega, \end{cases} \quad (4)$$

along a path of controls $u(x, \gamma)$ depending continuously on a parameter $\gamma \in [0, 1]$.

Despacito!

¹J.-M. Coron, E. Trélat, SIAM J. Control Optim. 43 (2004), no. 2, 549–569.   

Conclusions

- 1 We cannot get to the target 0 with non-negative controls,
- 2 But, we can, if the target, the objective y^1 is a positive steady state solution satisfying

$$\begin{cases} -\Delta y^1 = u^1 1_\omega & \text{in } \Omega \\ y^1 = 0 & \text{on } \partial\Omega, \end{cases} \quad (5)$$

with

$$u^1 \geq \nu > 0.$$

And the time T is long enough!!!

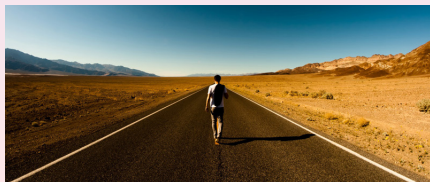


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Boundary control

The same results hold in the context of boundary control. In other words,

Theorem

Given a target y^1 , steady state solution (harmonic function) with boundary control $u^1 \geq \nu > 0$, the heat equation can be driven to it in time T large enough with control $u \geq 0$.

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We have shown that controllability with constraints can be achieved in long time. In fact, it is impossible to do it in short time!

Theorem

Whatever the initial datum y^0 and the steady state target y^1 associated to $u^1 \geq \nu$ is, the minimal control time under the positivity constraint is positive:

$$T_{\min} > 0,$$

except in the trivial case where $y^0 \equiv y^1$.

Proof of the waiting time

To fix ideas and without loss of generality we assume that $y^0 \equiv 0$.

The target $y^1 > 0$.

But we want to show that, if $u \geq 0$, it can not be reached in time T too short.

By duality, if $y(T) = y^1$,

$$\langle y^1, \varphi^T \rangle + \int_0^T \int_{\partial\Omega} u \frac{\partial\varphi}{\partial n} d\sigma(x) dt = 0,$$

where

$$\begin{cases} \varphi_t + \Delta\varphi = 0 & \text{in } (0, T) \times \Omega \\ \varphi = 0 & \text{on } (0, T) \times \partial\Omega \\ \varphi(T, x) = \varphi^T(x). & \text{in } \Omega \end{cases}$$

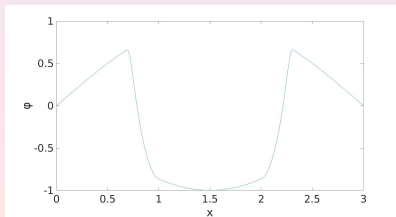
Thus, to conclude, in view of the identity

$$\langle y^1, \varphi^T \rangle + \int_0^T \int_{\partial\Omega} u \frac{\partial \varphi}{\partial n} d\sigma(x) dt = 0,$$

it suffices to find $T_0 > 0$ and a final datum $\varphi^T \in H_0^1(\Omega)$ such that, for any $T \in (0, T_0)$, the solution of the adjoint system with final datum φ^T satisfies:

$$\begin{cases} \left(\frac{\partial \varphi}{\partial n} \right)_+ = 0 & \text{on } (0, T_0) \times \partial\Omega \\ \langle y^1, \varphi^T \rangle < 0, & \forall T \in [0, T_0). \end{cases}$$

This is assured with an initial datum of the form



Explicit estimates on the waiting time

The proof above not only yields the fact that $T_{\min} > 0$, but actually gives lower bounds on this waiting time, by a careful analysis of the behaviour of the adjoint solutions.

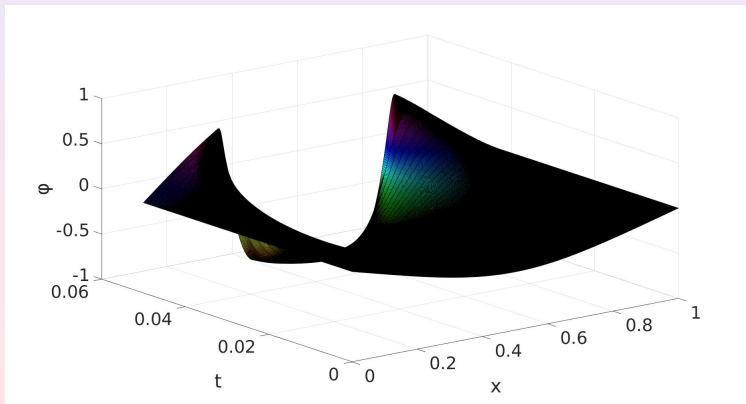


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Controllability in minimal time

Theorem

The system is controllable in minimal time T_{\min} with a non-negative measure u as control

The proof is again a consequence of the identity:

$$\langle y^1, \varphi^T \rangle + \int_0^T \int_{\partial\Omega} u \frac{\partial \varphi}{\partial n} d\sigma(x) dt = 0,$$

now applied to the solution of the adjoint heat equation φ with Φ_1 , the first eigenfunction of the Laplacian, as datum at time T :

$$\varphi = \exp(-\lambda_1(T - t))\Phi_1(x).$$

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The numerical method

- 1 We adopt a discrete approach: Discretise and minimise.
- 2 We discretise the equation both in space and time, with Δx and Δt small enough. We fix a number of time steps large enough, together with all the constraints (positivity of the boundary values, initial and final conditions).
- 3 We then minimise Δt using the IPOPT software.
- 4 We put a (very large upper bound) M on the control:

$$0 \leq u \leq M.$$

Numerical experiments

$y^0 \equiv 5$, $y^1 \equiv 1$, analytic estimate on minimal time : $\underline{T} \geq 0.165297$

See

<http://cmc.deusto.es/ipopt-and-ampl-use-to-solve-time-optimal->

Numerical experiments

$y^0 \equiv 1$, $y^1 \equiv 5$, analytic estimate on minimal time : $\underline{T} \geq 0.023076$

See

<http://cmc.deusto.es/ipopt-and-ampl-use-to-solve-time-optimal->

Optimality structure: And open problem

The numerical simulations above show that, apparently, the constrained controls in minimal time have the following properties:

- 1 They are unique.
- 2 They are constituted by the union of arcs where the controls vanish together with impulsive controls of Dirac delta type with support on a sequence of time instances concentrating at the final time.

This raises an interesting open problem. In view of the structure of the moment problem, can one find a solution constituted by the accumulation of a countable number of Dirac masses such that

$$u(t) = \sum_{j \geq 1} m_j \delta_{t=\tau_j}, \quad \sum_{j \geq 1} m_j < \infty$$

with $m_j \geq 0$ and $0 < \tau_j < T$ accumulating at T as $j \rightarrow \infty$ and such that the following identity hold for all $p \geq 0$?

$$\frac{2y^1}{((2p+1)\pi)^2} - \frac{e^{-(2p+1)^2\pi^2 T}}{(2p+1)\pi} y_{2p+1}^0 = 2 \sum_{j \geq 1} m_j e^{-(2p+1)^2\pi^2(T-\tau_j)}.$$

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As observed by **M. Tucsnak**, the waiting time phenomenon for the control under nonnegativity state constraints of the linear heat equation is related to that on the control of the **viscous Hamilton-Jacobi** equations proved by **A. Porretta and E. Z. Ann. Inst. Henri Poincaré, Anal. Non Linéaire, 2012**. If y solves the heat equation and $y(t, x) \geq 0$, the logarithmic change of variable $z(t, x) = -\ln y(t, x)$ leads to the viscous Hamilton-Jacobi equation

$$\dot{z} - \Delta z + |\nabla z|^2 = 0 \quad (t > 0, x \in \Omega), \quad (6)$$

with the initial condition $z(0, x) = -\ln yr^0(x)$ and constant target state $z^1 = -\ln yr^1$.

The waiting time for the control of this viscous Hamilton-Jacobi equation was proved using the **barrier functions** by Lasry and Lions. This is directly connected with the need of a minimal time for the constrained controllability of the linear heat equation. Barrier functions are achieved through the same logarithmic change of variables out of the **first eigenfunction of the Laplacian**.

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Conclusions

We have seen in a number of examples that:

- 1 There is a waiting phenomenon for controllability under constraints.
- 2 Constrained controllability can be achieved in a long enough time by two different methods: dissipativity and/or step-wise.
- 3 There is a measure control in the minimal time.
- 4 The numerical simulations show that the minimal time control is sparse or impulsional, composed by a sequence of diminishing Dirac deltas.
- 5 There is a link with the waiting time phenomena for the viscous Hamilton-Jacobi equation.

Open problems

Fully understand the sparsity structure of controls in minimal time